

PARTIALLY COHERENT PHOTON BEAMS FROM STORAGE RING UNDULATORS



Johannes Bahrtdt, Helmholtz-Zentrum Berlin, Chicago, March 15th, 2019
Coherence in particle and photon beams: Past, Present, and Future

Undulator Workshop 1989 @ BESSY

HELMHOLTZ
RESEARCH FOR GRAND CHALLENGES

HZB Helmholtz
Zentrum Berlin



Workshop on
Insertion Devices for Circularly Polarized Light
BESSY, Berlin November 1st and 2nd, 1993
Programm

18.00 SUNDAY OCTOBER 31st
Reception at BESSY

MONDAY NOVEMBER 1st

Morning Session

- 9.00 K.J. Kim LBL An Overview of Insertion Devices with Variably Polarized Light
- 9.40 P. Elleaume ESRF Circular Polarization from Undulators and Wigglers. Application to the ESRF
- 10.20 *coffee*
- 10.40 B. Diviacco ELETTRA Insertion Device Sources of Circularly Polarized Radiation at ELETTRA
- 11.20 S. Sasaki JAERI (topic to be announced)

Afternoon Session

- 14.00 A. Friedman NSLS Expected Performance of the AC Elliptical Wiggler
- 14.40 B. Kincaid ALS (topic to be announced)
- 15.20 *coffee and discussion*

TUESDAY NOVEMBER 2nd

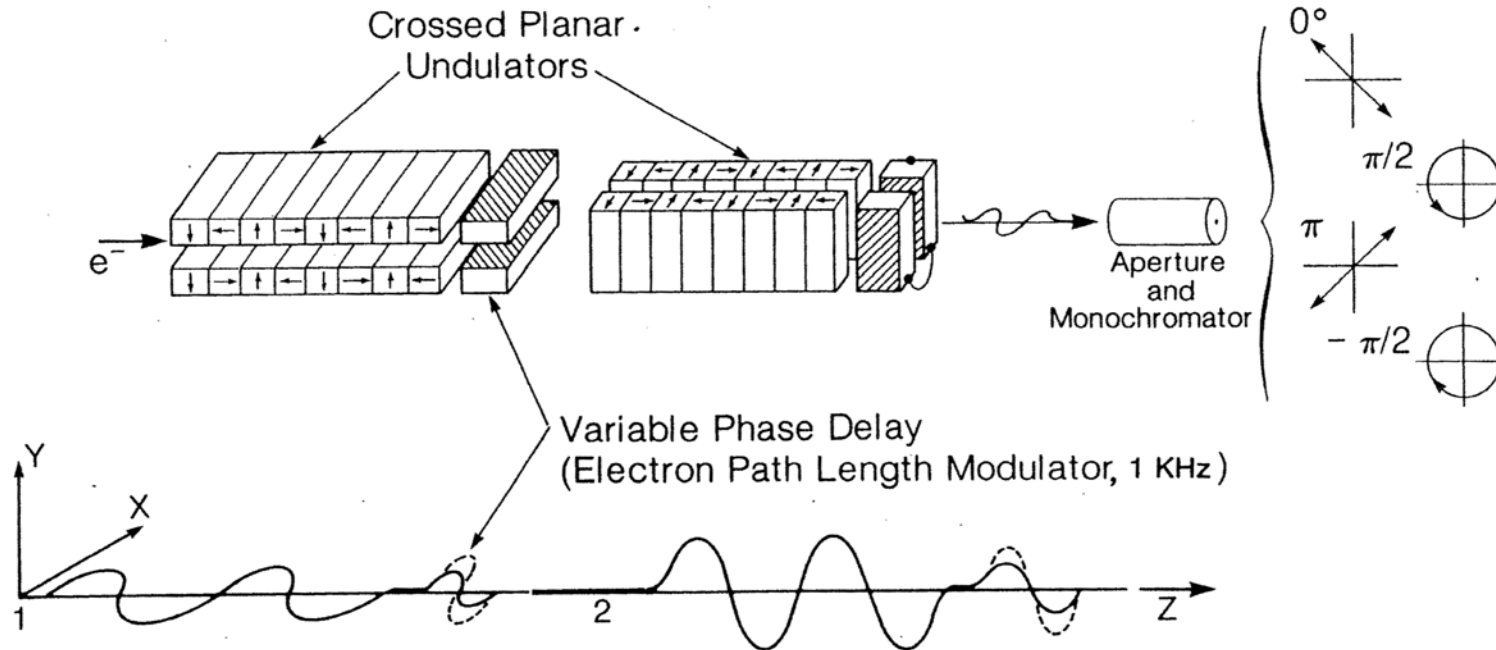
Morning Session

- 9.00 J. Pflüger HASYLAB The Asymmetric Wigglers at
- 9.40 B. Craft CAMD The Design of an Asymmetric and its Beamline Interface





Variably Polarized Radiation can be Generated with Crossed Undulators in Low Emittance Storage Rings



no longitudinal motion
fast polarization switching
(electromagnetic)

BESSY

XBL 8311-4539-A

Polarization change of undulator radiation

Moiseev, Nikitin, Fedosov

Russian

Izvestiya Vysshikh Uchebnykh Zavedenij, Fizika; (no.3); p. 76-80, 1978

English translation

Russian Physics Journal, Springer, (former Sov. Phys. J.)

21, 3 (1978) 332-335

Ideal parameters, no limiting factors discussed such as

- emittance
- energy spread
- beamline acceptance
- length of modules (# of periods)

Kwang-Je Kim, NIM **219** (1984) 425-429

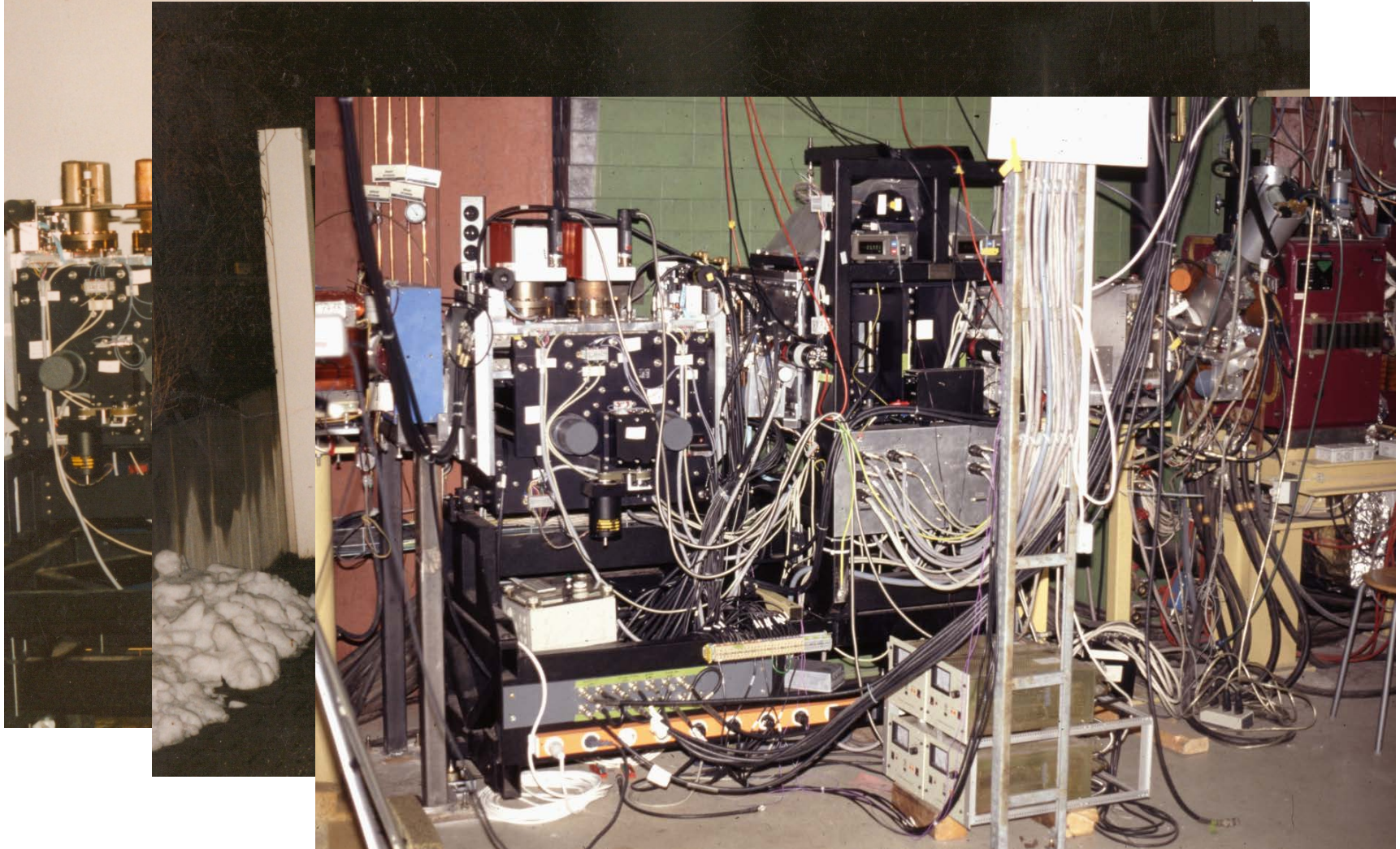
Evaluation of effects in old and new rings including

- emittance
- energy spread

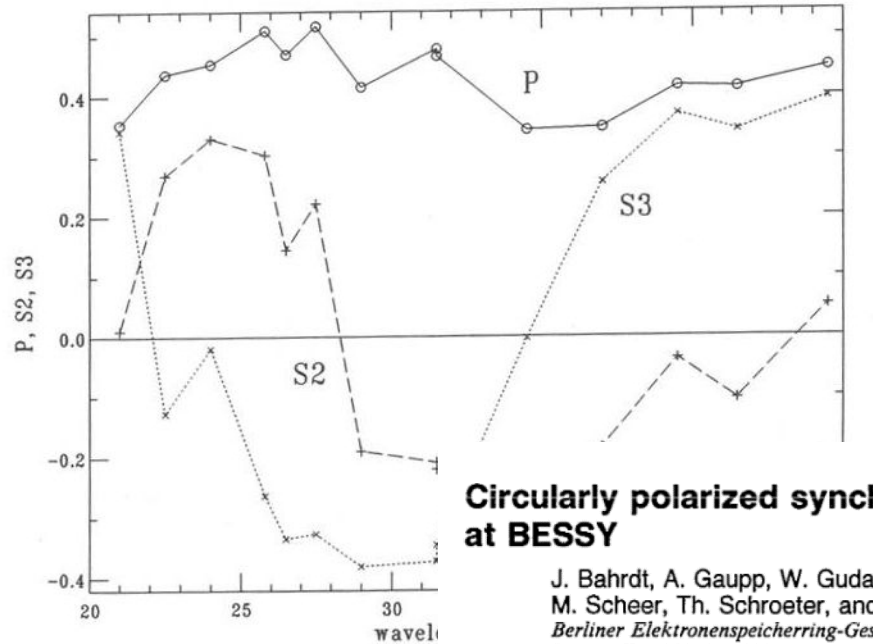
The BESSY I Crossed Undulator

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HZB Helmholtz
Zentrum Berlin



monochromator



Circularly polarized synchrotron radiation from the crossed undulator at BESSY

J. Bahrtdt, A. Gaupp, W. Gudat, M. Mast, K. Molter, W. B. Peatman, M. Scheer, Th. Schroeter, and Ch. Wang
Berliner Elektronenspeicherring-Gesellschaft für Synchrotronstrahlung m.b.H. (BESSY), Lentzeallee 100, D-1000 Berlin, Germany 33

(Presented on 15 July 1991)

The first experimental results from a double undulator producing circularly polarized synchrotron radiation are presented. The observed variation of the polarization parameters with photon energy is discussed. A strong dependence on the exact details of the tuning of the two undulators is observed. This probably accounts for the measured variation of the polarization parameters.

Post deadline Paper
SRI 1991, Chester

polarimeter

Onuki undulator:

- 1986 proposal
- 1989 realization

Period length

80mm

Number of periods

4

Magnetic gap

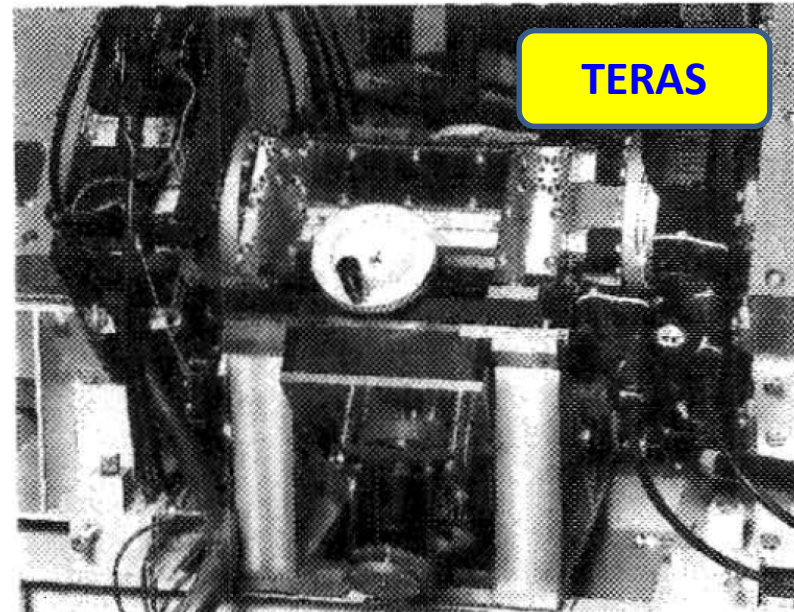
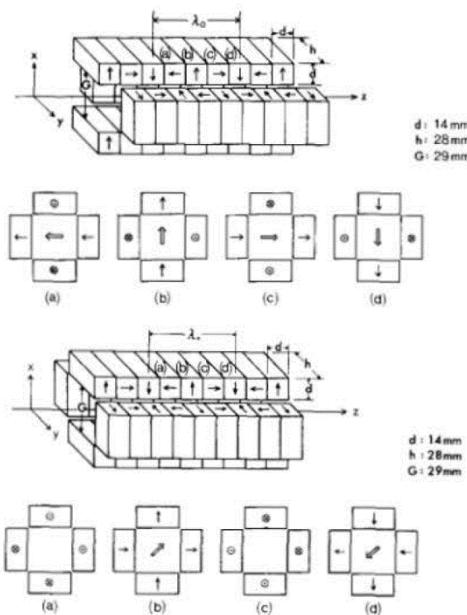
57 – 150mm

Polarization switching

3Hz

Design

permanent magnet



H. Onuki, NIM A246 (1986) 94-98

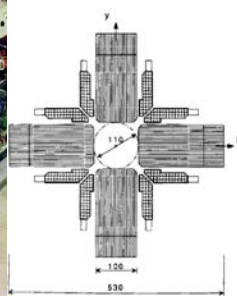
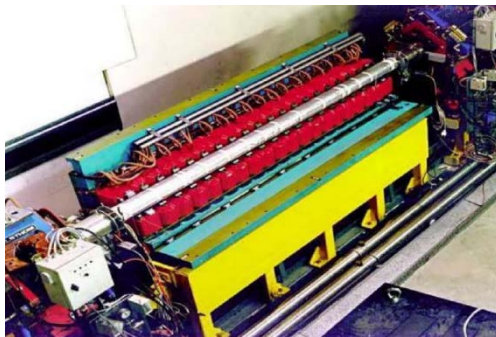
H. Onuki et al., RSI 60 (1989) 1838-1841

Onuki undulator:	Period length	86mm
1996 realization at NIJI-II	Number of periods	15
	Magnetic gap	64 – 160mm
	Polarization switching	3Hz
	Design	permanent magnet

T. Yamada, M. Yuri, H. Onuki, S. Ishizakaa, Rev. Sci. Instrum. 66 (2) (1995) 1493-1495

M. Yuri, K. Yagi, T. Yamada, H. Onuki, J. Electron Spectrosc. Relat. Phenom. (1996) 425-428

OPHELIE @ Super-ACO

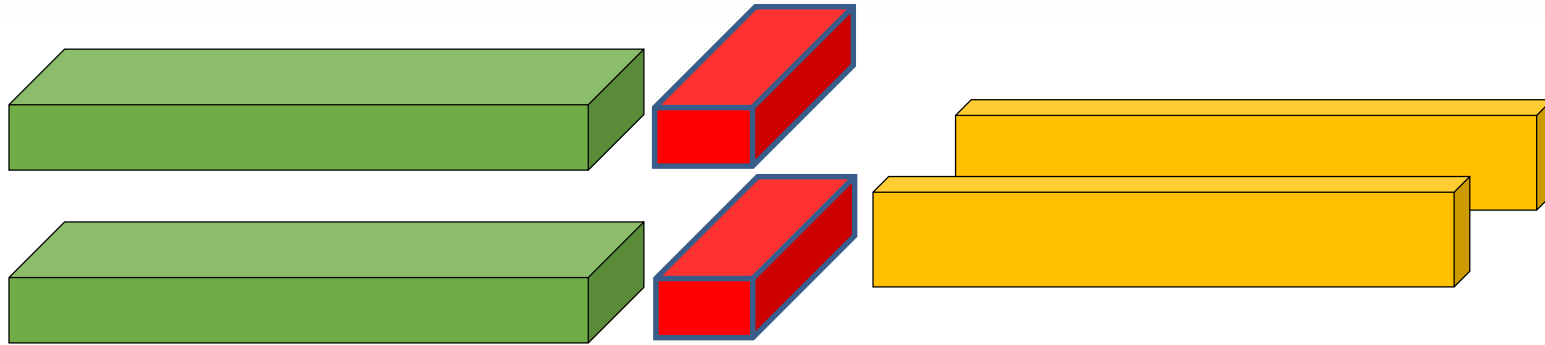


Period length	250mm
Number of periods	10
Magnetic gap	110mm
Polarization switching	1Hz
Design	electromagnet

L. Nahon et al., J. Synchrotron Rad. 5 (1998) 428-430

L. Nahon et al., NIM A 447 (2000) 569-586

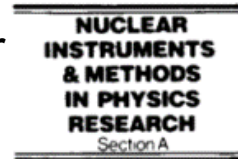
Crossed Undulator – Onuki Undulator – APPLE II



Onuki-Undulator (tilted by 45°) Phase shifter

Crossed undulator

Nuclear Instruments and Methods in Physics Research A 347 (1994) 83–86
North-Holland



Analyses for a planar variably-polarizing undulator

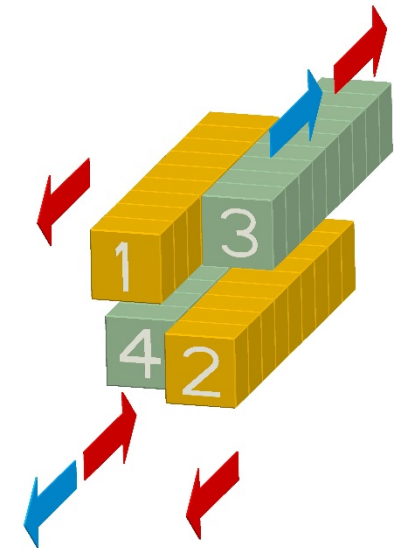
Shigemi Sasaki

Department of Synchrotron
Ibaraki 319-11, Japan

Institute, Tokai-mura, Naka-gun,

APPLE I 1993
APPLE II 1994

A new undulator for generation of synchrotron radiation is described in this paper. This device consists of two pairs of planar permanent magnet rows above and below the electron orbit plane. The magnetic field generated with this undulator, stronger than any other existing planar helical device, induces various types of electron motion such as vertically or horizontally sinusoidal motion and helical motion. The analyses of magnetic field and undulator radiation spectra are made on this undulator. For calculating the spectra and angular distributions in various modes, the SPring-8 storage ring parameters are assumed. This undulator generates brilliant circularly-polarized undulator radiation comparable in intensity with linearly-polarized radiation from a conventional undulator.



APPLE II: workhorse
at most 3rd generation SR



$$P = 1 - \frac{\langle(\Delta\alpha)^2\rangle}{2}$$

$$(\Delta\alpha)^2 = \left(\alpha_0 \frac{\sigma_\lambda}{\lambda}\right)^2 + 2 \left(2\pi\eta \frac{\gamma^2 \sigma_\theta^2}{1+K^2/2}\right)^2$$

energy spread & monochromator

emittance

original paper of Kim

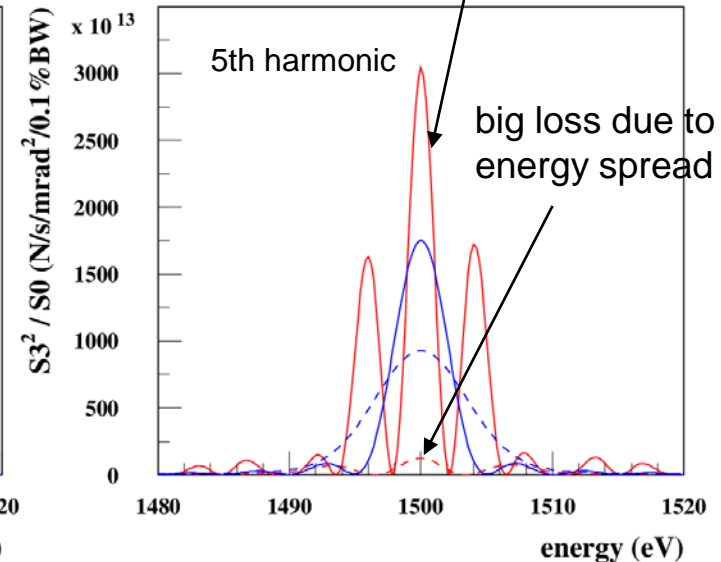
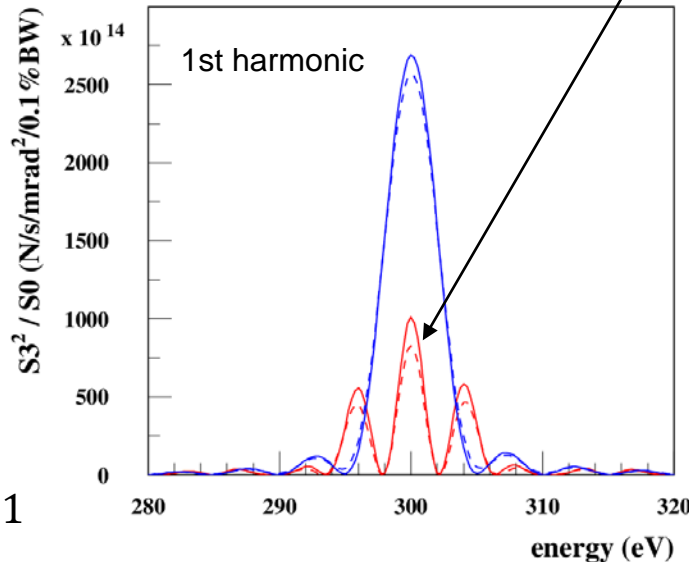
lower figure of merit due to less periods

larger figure of merit due to big S3 (helical)

$E = 1.72\text{GeV}$
 $\lambda_0 = 56\text{mm}$
 $\text{emittance} = 0$

solid: $\sigma_e = 0.0$

dashed: $\sigma_e = 0.001$



Crossed undulator: 2 x 30 periods
APPLE II: 1 x 60 periods

Segmented Undulator: Enhancement of Polarization Degree

- Increasing the degree of polarization with more segments

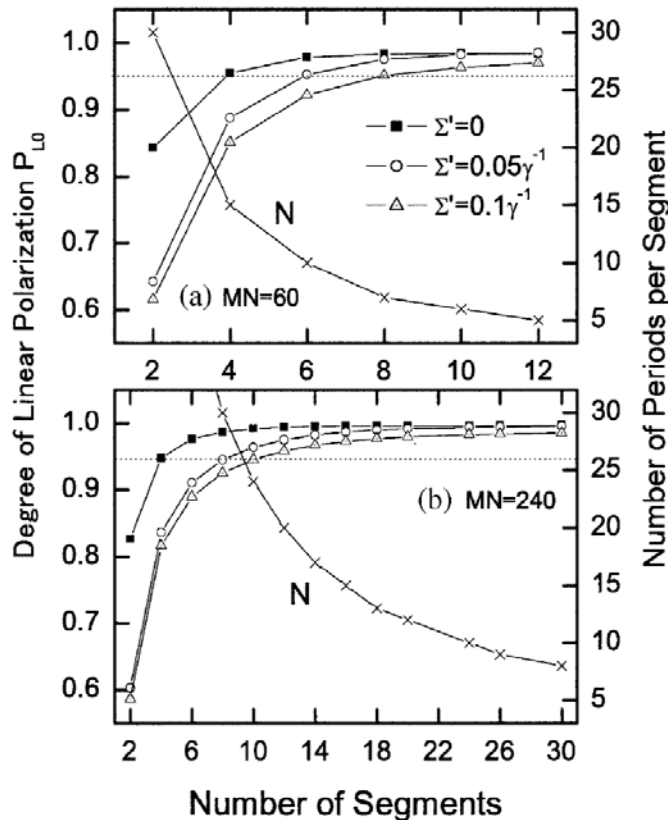


T. Tanaka, H. Kitamura, NIM A490 (2002) 583-591

$$I = 4 \left(\frac{2\pi MN}{\omega_1} \right)^2 S_N S_M$$

$$S_N = \frac{\sin^2[\pi N(1 - \omega/\omega_1)]}{\pi^2 N^2 (1 - \omega/\omega_1)^2}$$

$$S_M = \frac{\sin^2[(M/2)(2\pi N + \Phi_1)\omega/\omega_1]}{M^2 \sin^2[(2\pi N + \Phi_1)\omega/\omega_1]}$$



similarly: $\uparrow + \rightarrow \rightarrow \textcircled{\rightarrow} \textcircled{\rightarrow}$

T. Tanaka, H. Kitamura, SRI 2003, 231-234

Higher order suppression via
relative detuning of segments

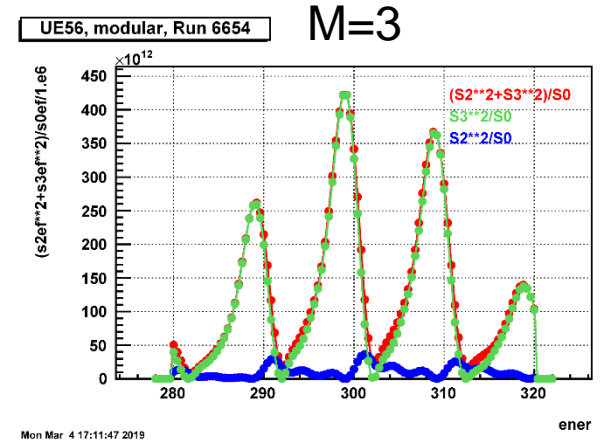
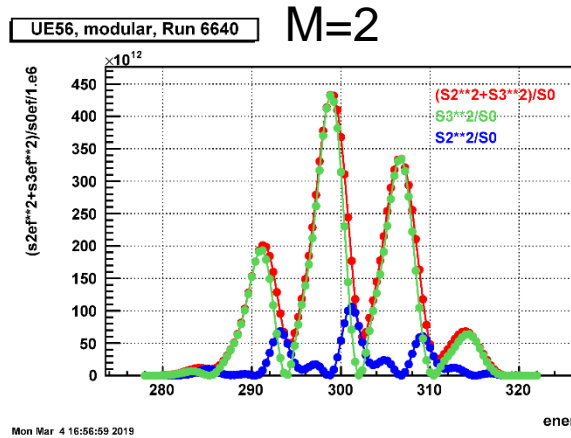
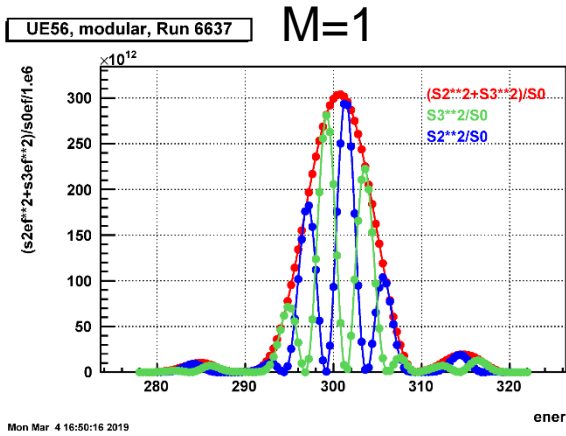
T. Tanaka, H. Kitamura, JSR 9 (2002) 166-269

Segmented Crossed Undulator Idea of Tanaka & Kitamura, 2002

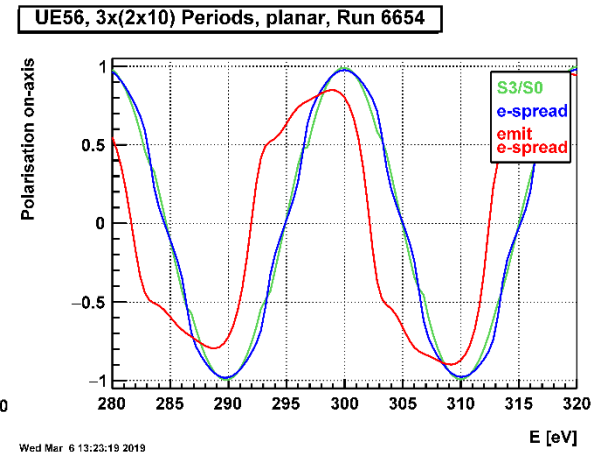
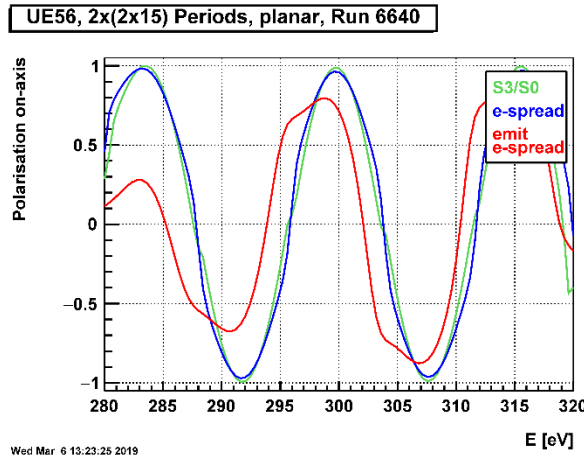
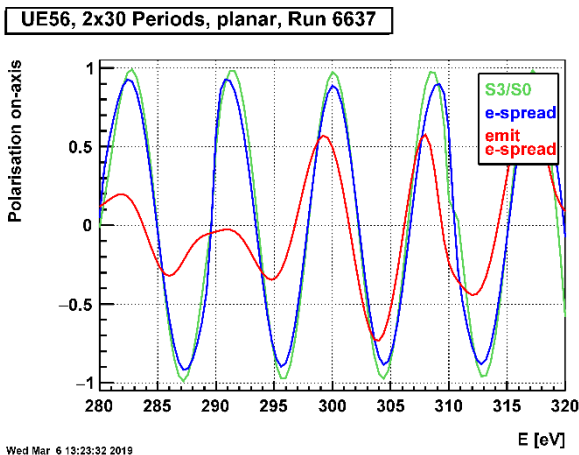
Periode 56mm
M=1, 2, 3
N=30, 15, 10
MN=30

Simulations with WAVE, Michael Scheer, HZB

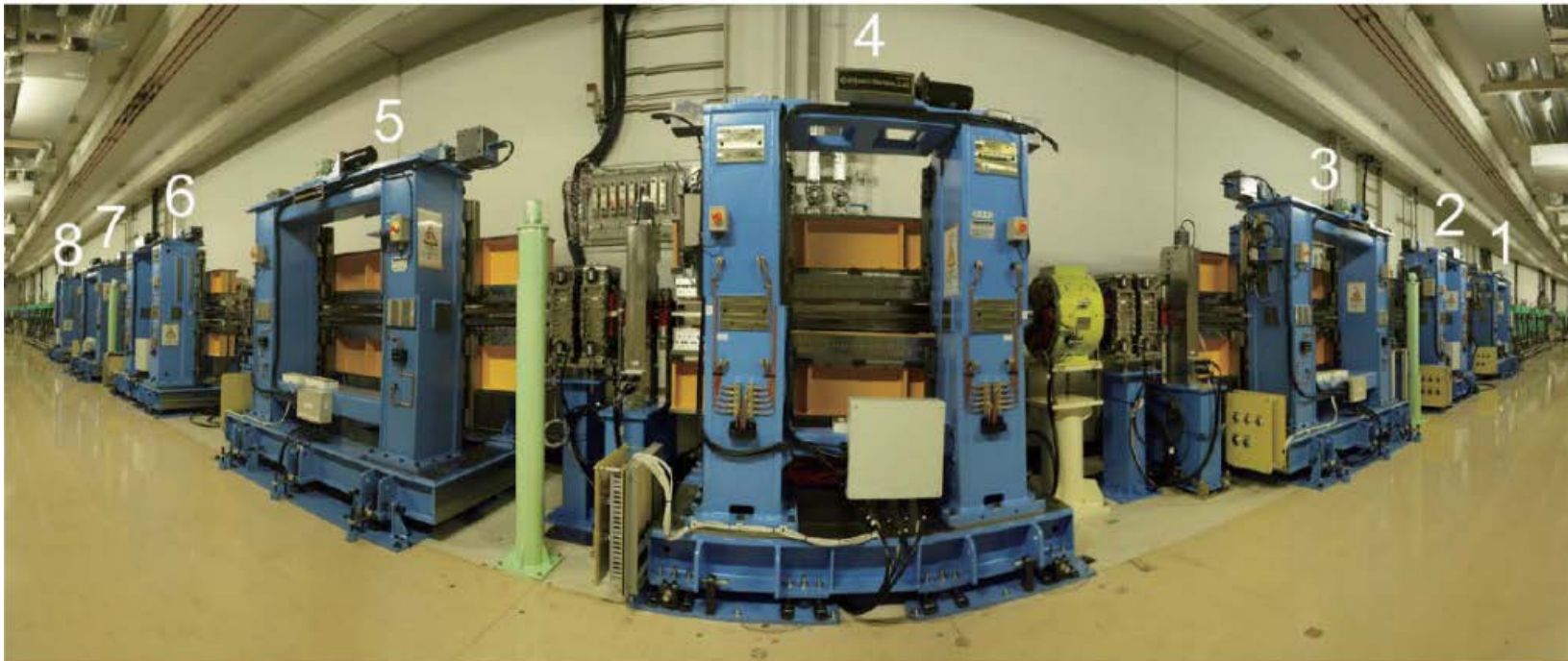
Flux density



Polarization degree



Segmented CU with low on-axis power: 8 Figure-8 Undulators

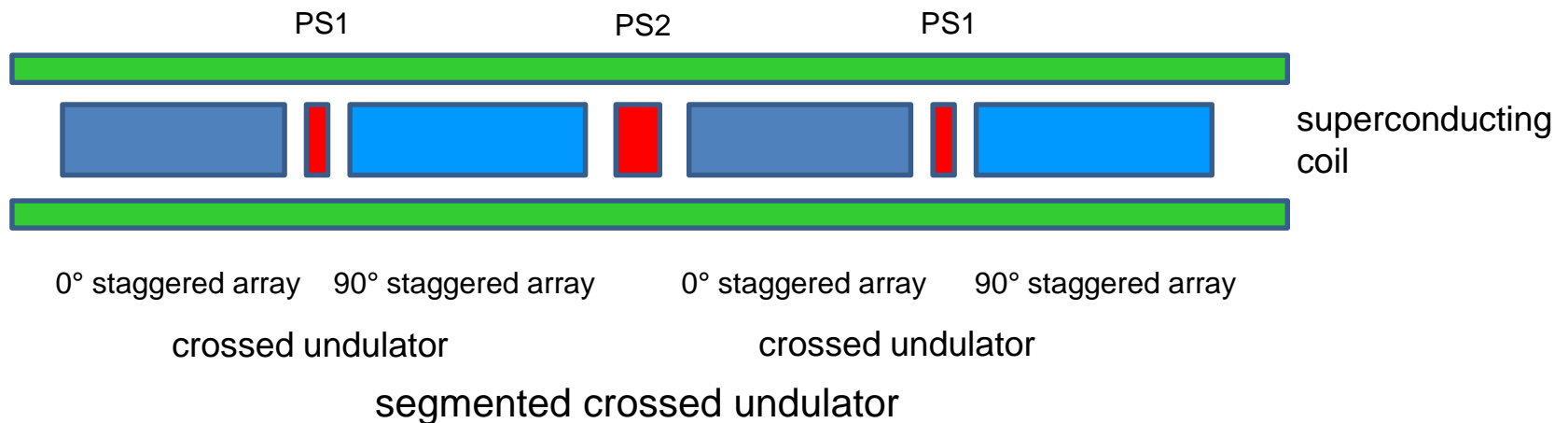


S. Yamamoto, *J. Synch. Rad.* (2014). 21, 352–365

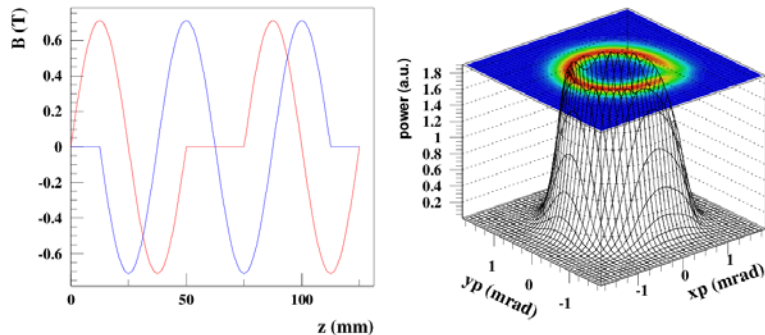
Highest fields at short periods
combined with variable polarization



HTS based staggered array
segmented crossed undulator

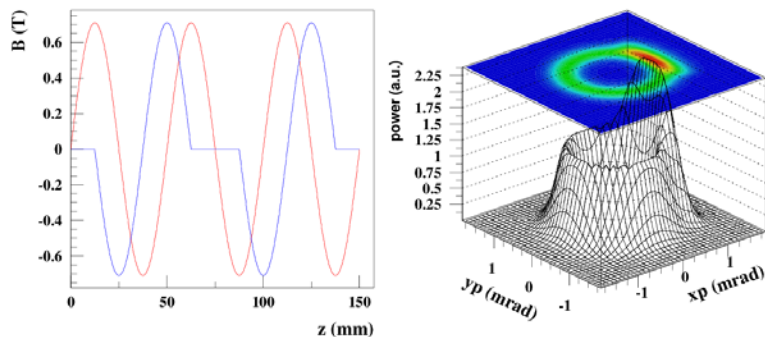


- High degree of polarization
- Linear polarized light without higher harmonics
- Linear polarized light with reduced on-axis power
- Even variable polarization is possible



Leaf undulator

J. Yan, S. Qiao, *Rev. Sci. Instrum.*, **81** (2010) 056101-1-3
S. Sasaki, *IPAC Taipeh, Taiwan 2018*
(concept for variable polarization)



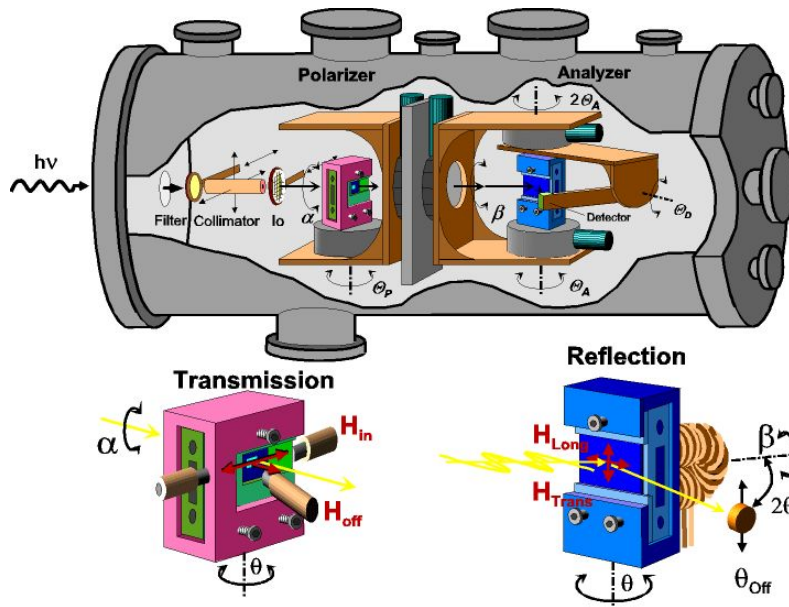
Knot undulator

S. Qiao et al., *Rev. Sci. Instrum.*, **80** (2009) 085108-1-4

APPLE knot undulator

S. Sasaki et al., *NA-PAC*, (2013) 1043-1045 (theory)
J. Fuhao et al., *JSR*, **22** (2015) 901-907 (built at SSRF)
Q. Zhou et al., *IEEE TRANS. APPL. SUPERCOND.*,
26,

(2016) 4101704-1-4 (built at SSRF)

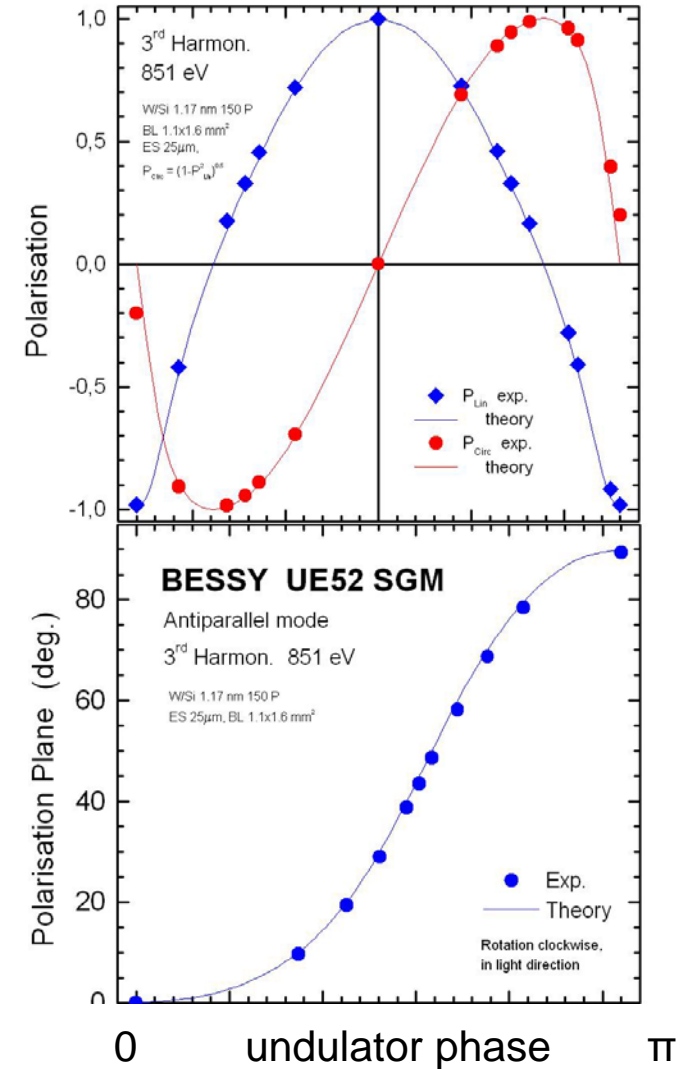


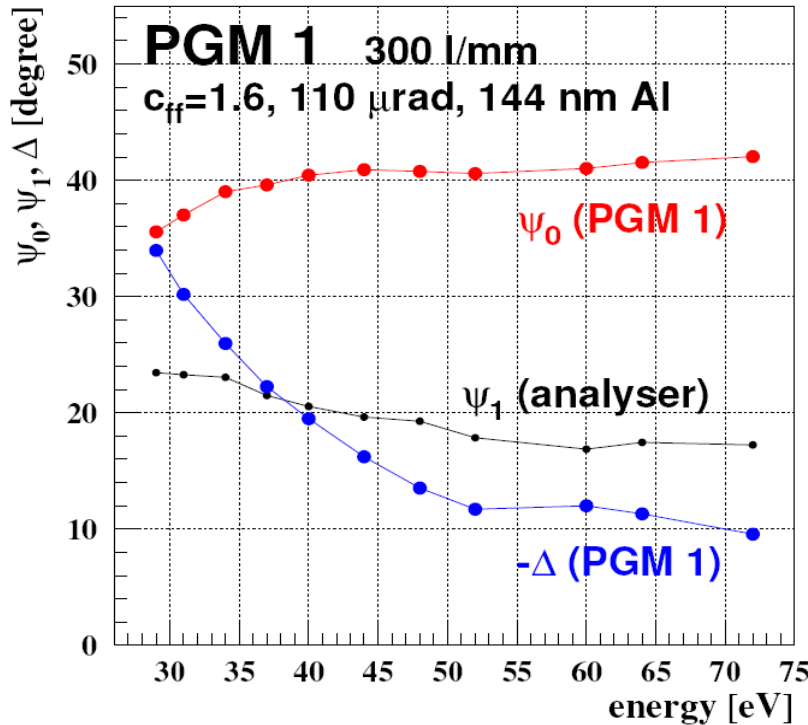
elliptical mode

inclined mode

HZB Soft X-ray Polarimeter
Excellent agreement between theory and measurement

F. Schaefers et al., Applied Optics 38 (1999) 4974





UE-112 APPLE @ BESSY II

$$\tan(\Psi_0) = |T_y^{BL} / T_x^{BL}|$$

transmissions of field amplitudes

$$\Delta = \phi_y^{BL} - \phi_x^{BL}$$

phase difference of electric field components

APPLE II produces any kind of polarization ellipse including a tilt in the so-called „universal mode“



circular polarization at the experiment

but also cases where polarization degree is spoilt in beamline



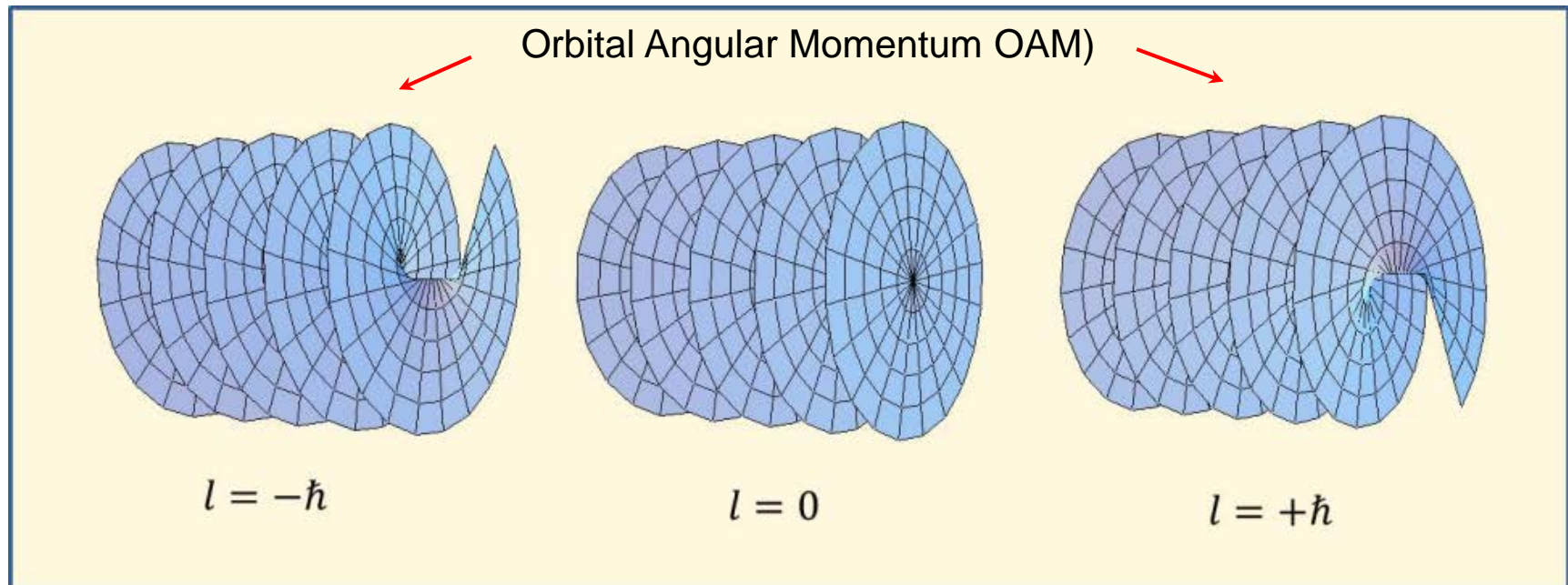
no recovery possible

J. Bahrtdt et al, Proc. PAC, Vancouver BC, USA (2009) 1162-1164

J. Bahrtdt et al, Conf. Proc. 1234 (2010) 335-338

Off axis radiation of higher harmonics of a helical undulator

$$A \sim \frac{(A_x - iA_y)}{\sqrt{2}} = \sqrt{2} e^{\pm i(n-1)\varphi} \left\{ \left(\gamma\theta - \frac{nK}{X} \right) J_n(X) - K J_n'(X) \right\}$$



S. Sasaki, I. McNulty and R. Dejus, NIM 582 A, 1 (2007) 43-46

S. Sasaki, I. McNulty, PRL 100 (2008) 124801

helical

$$A(r, \varphi) = \frac{a(r)}{L + d} \cos \left(\frac{\pi d}{\gamma^2 \lambda} + \frac{\pi}{(L + d)\lambda} r^2 \pm (n - 1)\varphi + \frac{2\pi L}{\lambda} - \omega t \right)$$

planar

$$B(r, \varphi) = \frac{b(r)}{L} \cos \left(\frac{\pi}{L\lambda} r^2 + \frac{2\pi L}{\lambda} - \omega t \right)$$

$$I(r, \varphi) = \frac{\omega}{2\pi} \int_0^{2\pi} (A + B)^2 dt =$$

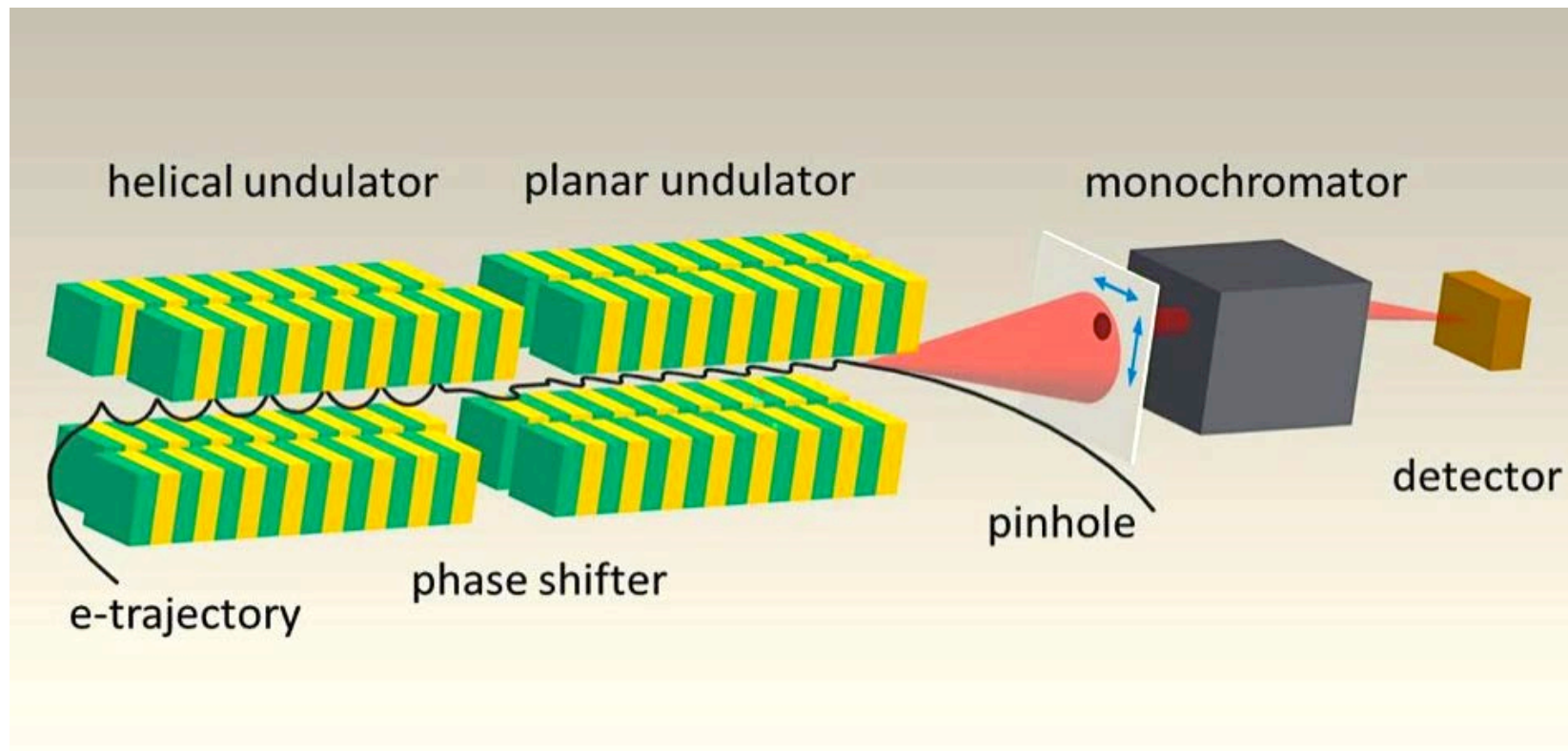
$$\frac{a^2}{2(L + d)^2} + \frac{b^2}{2L^2} + \frac{ab}{L(L + d)} \cos \left(\frac{\pi d}{\gamma^2 \lambda} - \frac{\pi d}{L^2 \lambda} r^2 \pm (n - 1)\varphi \right)$$

$$\varphi = \pm \left(-\frac{\pi d}{\gamma^2 \lambda} + \frac{\pi d}{L^2 \lambda} r^2 \right) / (n - 1)$$

This spiral structure of the intensity can be measured

Double APPLE II Undulator UE56
pinhole in 13mm distance, pinhole scans

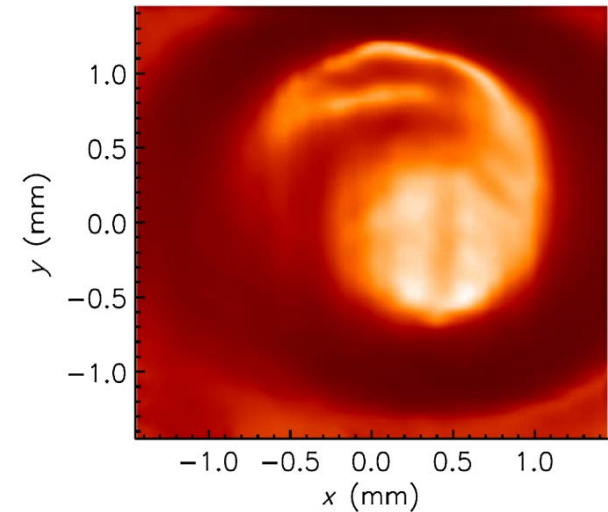
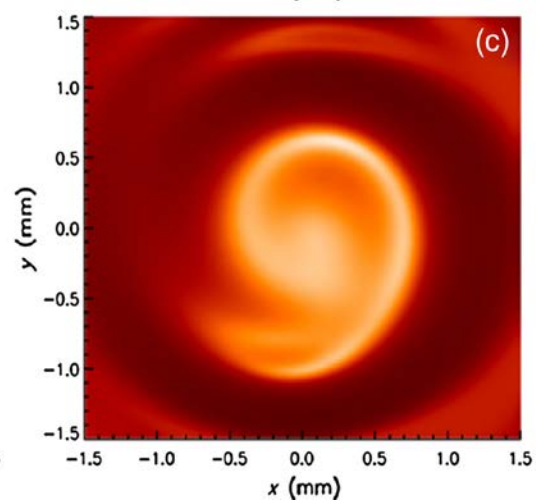
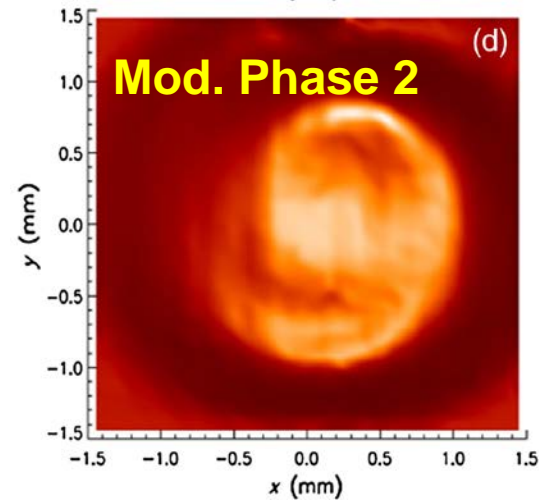
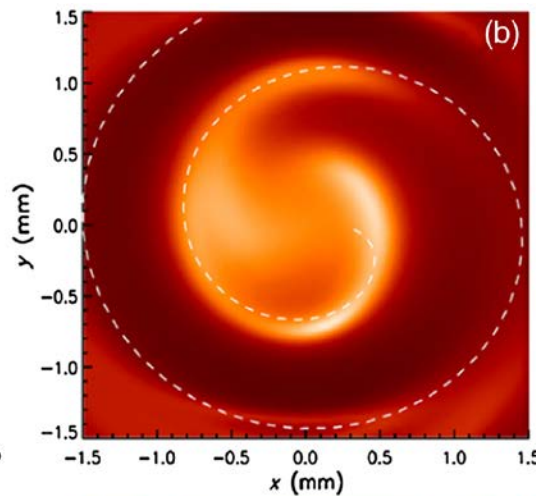
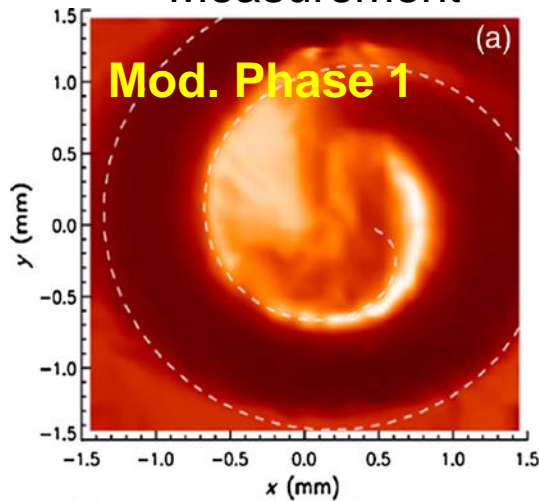
Spectral and spatial overlap of
2nd harmonic of helical ID & 1st harmonic of planar ID



J. Bahrtdt et al., Physical Review Letters 111 (2013), p. 034801/1-5

Measurement

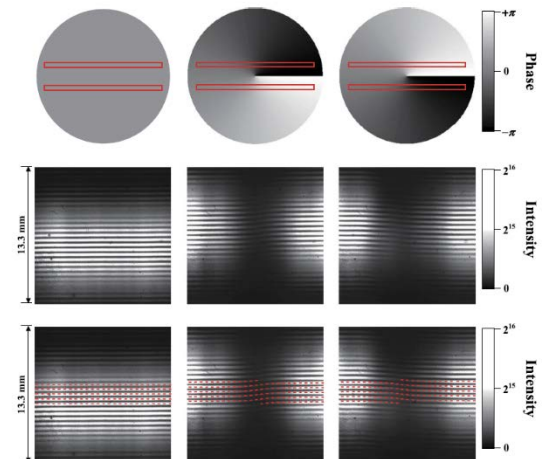
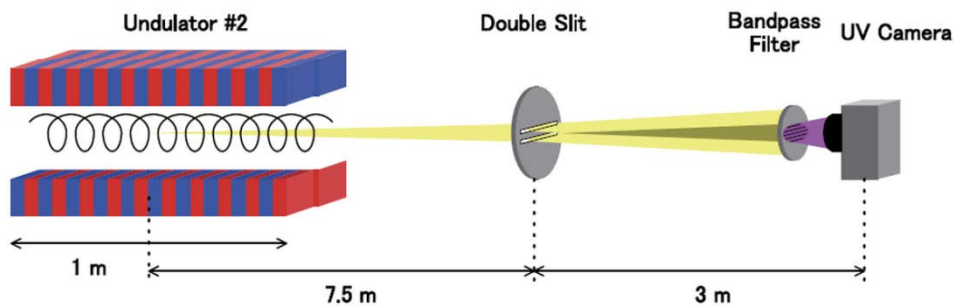
Simulation



dashed line depicts equal phase

$$\varphi = \pm \left(-\frac{\pi d}{\gamma^2 \lambda} + \frac{\pi d}{L^2 \lambda} r^2 \right) / (n - 1)$$

- Not successful at BESSY UE112 due to optical element quality
- Successful at UVSOR because:
 - larger wavelength: 355nm instead of 12.5nm at BESSY
 - Energy selection with band pass filter instead of monochromator (no degradation due to optical element quality)



Katoh et al., Scientific Reports, 7: 6130 (2017) 1-8

Generation of photon beams with high degree of coherence (ID-shimming)
minimization of electron beam effects (emittance reduction)

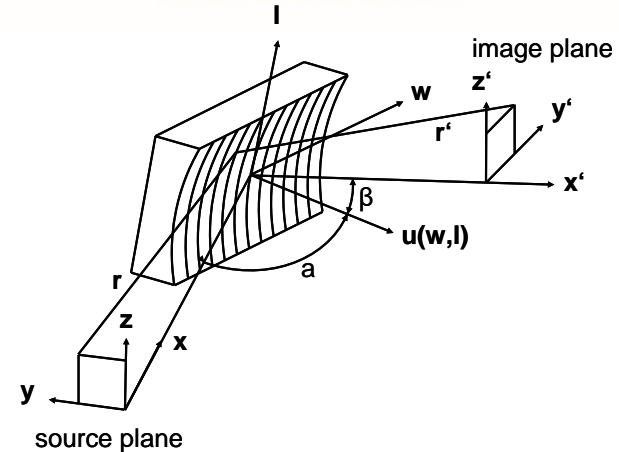
Keeping the coherence during propagation to the experiment
pushing the optical element quality further

Methods of partially coherent photon beam propagation:

- Fourier Optics (FO), linear systems
J.W.Goodman, Introduction to Fourier Optics, McGRAWHILL, 1968
- FO + local ray-tracing including the phase, SRW
O. Chubar et al., J. of Physics: Conf. Ser. 425 (2013) 162001
- Wigner function (brightness) of real sources
G. Geloni, V. Kocharyan, E. Saldin, JSR, 22 (2015) 288-316
- Propagation of Wigner function, SPECTRA
T. Tanaka, PRST-AB, 17 (2014) 060702-1-14
- Stationary phase approximation, PHASE
J. Bahrtdt, PRST-AB, 10 (2007) 060701-1-15
J. Bahrtdt, U. Flechsig, S. Gerhardt und I. Schneider Proc. of SPIE 8141, Advances in Computational Methods for X-Ray Optics II, Bd. 8141, 2011.
Source code on github: <https://github.com/flechsig/phase>

Propagation of electric fields
along a single optical element is
based on power series expansions of

- coordinate tranformation
- path length
- determinants
- ...



$$E(y', z') = \frac{\sqrt{\cos(\alpha)} \sqrt{\cos(\beta)}}{\lambda^2} \iint E(y, z) \left\{ \iint \frac{1}{r \cdot r'} e^{ikPL} dw \cdot dl \right\} \left| \frac{\partial(y, z)}{\partial(dy', dz')} \right| ddy' \cdot ddz'$$

2nd order expansion of path length (PL)

$$PL(w_0 + \Delta w, l_0 + \Delta l) = PL_{w_0 l_0} + \frac{1}{2} \frac{\partial^2 PL}{\partial w^2} \Bigg|_{w_0 l_0} \Delta w^2 + \frac{1}{2} \frac{\partial^2 PL}{\partial l^2} \Bigg|_{w_0 l_0} \Delta l^2 + \frac{\partial^2 PL}{\partial w \cdot \partial l} \Bigg|_{w_0 l_0} \Delta w \cdot \Delta l$$

and analytic integration over the optical element surface;

for implementation of emittance & energy spread: integration over phase space

3rd order derivative of PL needs to be implemented



high brightness sunshine
beautiful non-linear mountain trails
and always great views